



## Ultra Low Noise CMOS preamplifier-shaper for X-ray spectroscopy

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### Abstract

We report here an Ultra Low Noise CMOS charge-sensitive amplifier (CSA) and a CSA-Shaper for X-ray spectroscopy, suitable for a small anode capacitance, low leakage current solid state detector. The continuously sensitive CSA and CSA-Shaper use a MOS transistor MF to discharge the integration capacitance; a novel self-bias circuit biases MF in the  $G\Omega$  region (subthreshold region) while tracking process, temperature and supply voltage variations. A T-Network in the DC-feedback improves the CSA gain linearity (within 0.1% up to 1.8 fC input charge).

The CSA-Shaper features a novel Adaptive Pole-Zero cancellation scheme: the Zero in the Shaper adapts itself statically and dynamically to cancel the CSA Pole.

The noise performance of the CSA has been measured: with no detector connected and MF pinched, the room-temperature equivalent noise charge (ENC) is 9e rms at 12  $\mu$ s shaping time. When connected to a cooled detector ( $T = -75^\circ\text{C}$ ), with MF in pinch off, an ENC of 13e rms is achieved (FWHM = 111 eV) at 2.4  $\mu$ s shaping time. This is the best reported energy resolution ever obtained with a CMOS preamplifier. © 1998 Elsevier Science B.V. All rights reserved.

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### 1. Introduction

For ultra low noise applications, front-end circuits matched to detectors with small anode capacitance and low leakage current are used. The charge-sensitive amplifier (CSA) is widely used, as its voltage-charge gain is independent of anode capacitance variations (Fig. 1). The signal charge is integrated on a capacitance CF, which has to be discharged to prevent the preamplifier from being saturated by subsequent integration events and the

detector leakage current. The technique used to discharge CF influences significantly on the performance of ultra low noise systems [1,2].

Three approaches have been proposed so far for the front-end circuit. The first one is the integration of the input transistor of the preamplifier into the detector itself to avoid the capacitance associated with the connections [3]. This approach gave the best performance so far [4], but its drawback is the fabrication of the CSA's input transistor with all the constraints imposed by the production of high-quality detectors. As a consequence: (i) the first transistor is not optimized, (ii) the introduction of additional masks reduces the yield of the process

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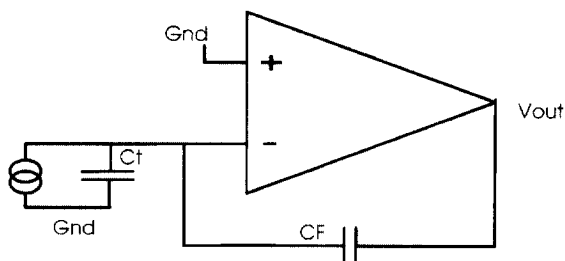


Fig. 1. Charge-Sensitive Amplifier (CSA) based on a core voltage amplifier with shunt feedback.

(iii), the integrated transistor increases the usually negligible anode capacitance.

The second approach is to use an external JFET as input transistor of the CSA. An X-ray system presently under test implements a novel feedback scheme by using a JFET transistor with input capacitance of 0.5 pF [5].

The third approach is the use of a monolithic CMOS preamplifier [6–8]. A *P*-channel input device is mandatory to benefit from the lower flicker noise (compared to a *N*-channel one). The use of low-cost CMOS technology calls for adhoc design techniques. In this paper we report novel circuit techniques and their implementation on a CMOS front-end circuit for ultra low noise applications.

## 2. Circuit techniques

To establish the DC path, a resistance  $R$  can be inserted in parallel to the integration capacitance  $C_F$ .  $R$  does not contribute to the parallel noise if  $4kT/R \ll 2qI_{leak}$ . Ultra low noise systems require detectors with leakage currents  $1 \text{ pA} < I_{leak} < 100 \text{ pA}$ , thus  $R$  must be in the range.  $5 \text{ G}\Omega < R < 50 \text{ G}\Omega$ . The only possible way to realize such a high resistance in CMOS technology is by using a MOS transistor. To fully benefit from the small anode capacitance, a minimum size feedback device with negligible parasitics added to the input node is required. These requirements set the biasing condition of MF in the subthreshold region, where its associated resistance  $R_{ds}(\text{MF})$  changes exponentially with the bias potentials.

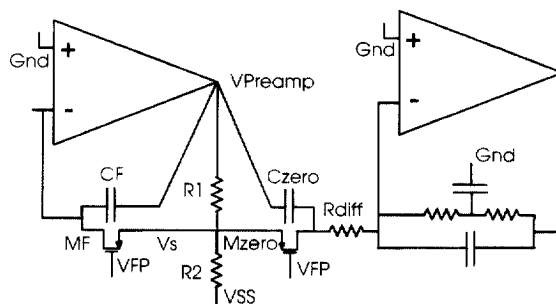


Fig. 2. CSA (left-hand side) and Shaper circuit.

The use of such device in the feedback loop necessitates the solution of three issues: (i) how to bias MF in the  $\text{G}\Omega$  range, and stabilize its DC resistance against process, temperature and supply voltage variations (in particular, the threshold voltages can change  $\pm 100 \text{ mV}$  from chip to chip); (ii) how to improve the linearity of a system with a non-linear resistance in the feedback loop. The worsening is due to the change of the decay time during the transient and its dependence upon the integrated charge. Finally (iii) how to stabilize the voltage-charge gain versus the detector leakage current.

We addressed those issues by using novel circuit techniques. The techniques have been implemented in a CSA [9] and a CSA-Shaper fabricated in two different technologies (HP  $1.2 \mu\text{m}$  and AMS  $1.2 \mu\text{m}$ ).

### 2.1. T-network in the DC feedback

In a real CSA, the signal charge is integrated onto  $C_F$  with a rise time  $T_{rise}$  and discharged with a decay time  $T_{decay}$  imposed by the feedback network. The voltage-charge gain, which in the ideal case is  $1/C_F$ , is reduced when the ratio  $T_{decay}/T_{rise}$  is not sufficiently high. In such conditions, a variation of  $T_{decay}$  results in a variation of the voltage-charge gain.

In our case  $R_{ds}(\text{MF})$  depends exponentially upon the biasing conditions of MF, thus its value changes during the discharge of  $C_F$ . If  $C_F = 15 \text{ fF}$ , with  $Q_{in} = 0.75 \text{ fC}$  the output swings of  $50 \text{ mV}$  which, if applied to the  $V_{gs}(\text{MF})$  would change

$R_{ds}(MF)$  of a factor four. To improve the linearity of the voltage–charge gain we inserted a voltage divider between the output of the preamplifier and the source of MF (Fig. 2). Consequently, the  $R_{ds}(MF)$  variation during the dynamic swing is attenuated. Moreover,  $T_{decay} = k \cdot R_{ds}(MF) \cdot CF$ , i.e. the ratio  $T_{decay}/T_{rise}$  is increased by a factor  $k = (R1 + R2)/R2$ .

We measured the voltage–charge gain using an on-chip injection capacitance and a Multichannel Analyzer. The CSA gain ( $G = 60 \text{ mV/fC}$ ) is linear within 0.1% up to 1.8 fC input charge.

## 2.2. Self-adaptive bias circuit for MF

We developed and implemented the Self-Adaptive bias scheme for a CMOS CSA with  $N$ -channel input device, matched to 0.8 pF detector capacitance, to be used for a particle tracking system [10,11]. In that application the peaking time was 50 ns, and we achieved a resolution response of 6 MΩ with a  $W/L = \frac{3}{20}$  MOS device biased in the triode strong inversion region.

In the present application the problem is similar, but the scheme must be modified for use with a  $P$ -channel input device. Moreover, the requirements are more challenging: we are forced to use a minimum size MOSFET device ( $W/L = 3 \mu\text{m}/3 \mu\text{m}$ ) biased in the subthreshold region, with an associated resistance three orders of magnitude larger (GΩ region). This makes  $R_{ds}(MF)$  extremely sensitive to the biasing condition of MF and on the threshold voltages. The issue has been addressed by modifying the self-adaptive scheme for a  $P$ -channel input device and improving its capability to track threshold variations. The circuit shown in Fig. 3 keeps  $V_{gs}(MF) - V_{th}(MF) = \text{const.}$ , even when  $V_{gs}(MF)$  is only a few times  $kT/q$  below threshold.

The gate voltage  $V_g(MF)$  tracks the variations in  $V_{th}(M1)$  and  $V_{th}(MF)$ : a diode-connected replica MN1 of M1 is biased at the same current density so that the voltage at the drain of MN1 tracks the NMOS threshold ( $V_g(MN1) = V_g(M1)$ ). Moreover, the required gate-to-source voltage of the feedback device is generated by passing a current  $I_d(MNF)$  through the diode MNF. MNF is a scaled copy of MF created by laying out  $n$  parallel copies of MF.

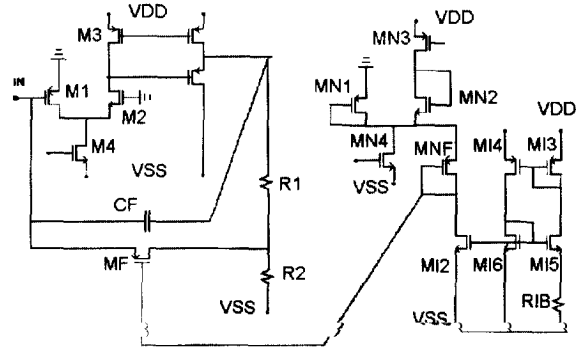


Fig. 3. CSA (left-hand side) and Self-Bias circuit for MF.

Since MNF has the same  $V_{bs}$  and width as MF, it experiences the same body and narrow-channel effects which affect  $V_{th}(MF)$ . It results that  $V_{th}(MF) = V_{th}(MNF)$  and

$$V_{gs}(MF) - V_{th}(MF) = V_{gs}(MNF) - V_{th}(MNF) \\ = F(I_{ds}(MNF)).$$

In other words, the biasing conditions of MF are set by the current  $I_{ds}(MNF)$  (which can be controlled within 20%) rather than by the threshold voltages, which influence  $R_{ds}(MF)$  exponentially.

In this circuit the number of MF's copies we use to lay-out MNF is  $n = 180$ . The low-current source (MI2–MI6, RIB in Fig. 3) set  $I_d(MNF) = 12 \text{ nA}$ .

In Fig. 4 the CSA response to charge pulse at different temperature is shown. In the 100°C temperature range  $V_{th}(MF)$  change of 50 mV and  $R_{ds}(MF)$  would change of two orders of magnitudes. The self-adaptive bias circuit tracks the threshold voltage and  $dR_{ds}(MF)/R_{ds}(MF) = 20\%$ .

The CSA output is moved positive by the detector leakage current  $I_{leak}$  and the node  $V_s$  (right-hand side of MF) becomes MF's effective source. The consequent reduction of  $R_{ds}(MF)$  (proportional to  $I_{leak}$ ) extends the range of the allowable  $I_{leak}$  current from the detector. Clearly a higher value of  $I_{leak}$  degrades the energy resolution.

It's possible to force on the node VFP (gate of MF) bias conditions different from the natural one. This feature was used to obtain the best possible noise performance.

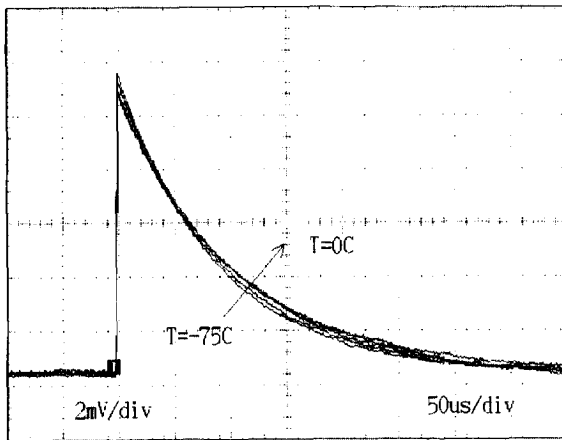


Fig. 4. Response of the CSA to a charge signal at different temperatures (–75 °C, –50 °C, –25 °C, 0 °C, 25 °C).

### 2.3. Adaptive Pole-Zero cancellation in the CSA-shaper

The new scheme, a variation of the one implemented in Refs. [9,10], is based on the classical Pole-Zero cancellation technique [12]. With reference to Fig. 2,  $M_{zero}$  is biased in the same way as MF for every value of  $I_{leak}$  and every value of  $V_{Preamp}$ . Consequently, the variations of  $R_{ds}(M_{zero})$  and  $R_{ds}(MF)$  during the discharge of CF are the same. In other words, the Zero associated with the

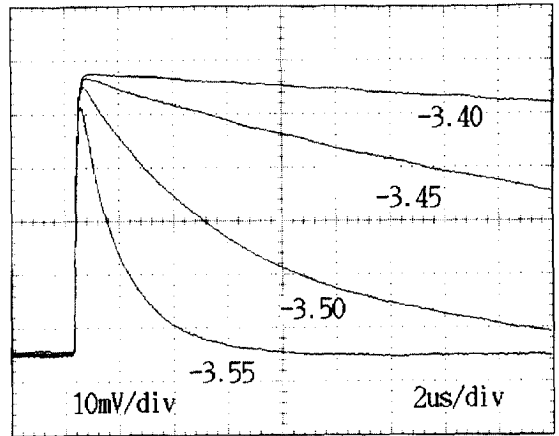
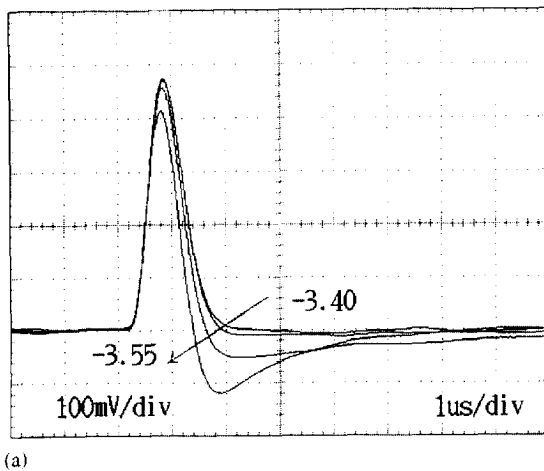
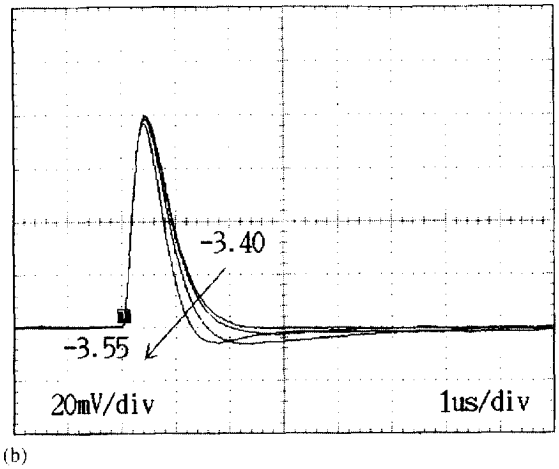


Fig. 5. CSA response to a charge signal for different  $V_{IP}$  forced to the gate of MF (–3.40, –3.45, –3.50, –3.55) V, i.e. feedback resistances.

network  $C_{zero}$ – $M_{zero}$  adapts itself *statically* and *dynamically* to accurately cancel the preamplifier pole. The response of the preamplifier to a charge event with various VFP forced to the gate of MF (i.e. various feedback resistances) is reported in Fig. 5. The corresponding responses of the CSA-External Shaper and CSA-Shaper (Fig. 6) prove the effectiveness of the scheme: the adaptive Pole-Zero cancellation makes the response of the CSA-Shaper independent on the position of the CSA Pole.



(a)



(b)

Fig. 6. (a) CSA-external Shaper and (b) CSA-Shaper response to a charge signal at various  $V_{IP}$  forced to the gate of MF and  $M_{zero}$ .

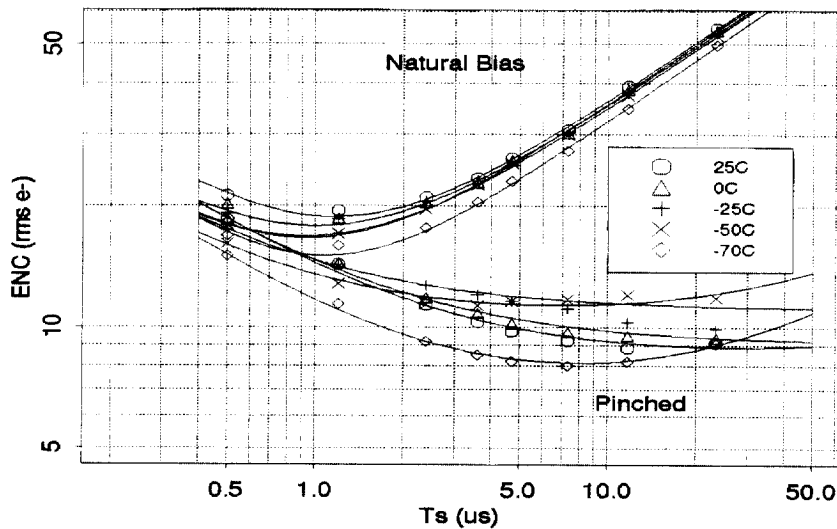


Fig. 7. ENC versus shaping time for the AMS CSA at various temperatures (CSA not connected to detector). MF self-biased (upper group) and MF pinched-off (lower diagrams).

### 3. Noise performance of the CSA

We measured the ENC versus shaping time for various temperatures in two bias conditions of MF: pinch-off, and the one imposed by the self-adaptive bias circuit ( $1.8 \text{ G}\Omega$ ). As far as the pinched state is concerned, an increase of the gate voltage of the feedback device MF by about  $15 \text{ mV}$  from the self-biased condition, brings MF's effective resistance to about  $100 \text{ G}\Omega$ . Under these conditions, the parallel noise is almost eliminated, as seen in Fig. 7. However, the circuit is still stable and continuously sensitive and can be maintained in this condition without adjustment of the gate voltage for up to  $12 \text{ h}$ ; this is sufficient for most spectroscopy runs. The only practical use of the CSA in this pinched condition is with detectors having  $I_{\text{leak}} < 1 \text{ pA}$ . With no detector connected, we measure an ENC of  $9 \text{ e rms}$  at room temperature.

The CSA, connected to an external Shaper, has been coupled to a Silicon Drift Detector (SDD) [13] cooled down to  $T = -75^\circ\text{C}$ . The capacitance of the SDD anode plus the parasitics associated with the bond wire are estimated equal to  $600 \text{ fF}$ . With  $t_s = 2.4 \mu\text{s}$  shaping time and MF in pinch-off, the spectrum of  $^{241}\text{Am}$  and  $^{55}\text{Fe}$  has been taken

(Figs. 8 and 9). We measure  $\text{FWHM} = 111 \text{ eV}$  ( $\text{ENC} = 13 \text{ e rms}$ ), which is an improvement of almost a factor of 2 over the best previously reported resolution of a Si detector coupled to a CMOS circuit [6,7]. The noise performance is summarized in Table 1.

These noise performances have been achieved with a pseudo-Gaussian Shaper with the CSA's noise dominated by the flicker noise of M1 (long shaping time, low  $I_{\text{leak}}$ ). An optimal filter for flicker noise improved the energy resolution by  $18\%$  [14,15]: we estimate that with such a filter our CSA connected to a detector can reach  $11 \text{ e rms}$ .

The noise performance of the CSA-Shaper was not measured.

### 4. Conclusions

Novel circuit techniques have been implemented in an ultra low noise CSA and a CSA-Shaper fabricated in two different technologies (HP  $1.2 \mu\text{m}$  and AMS  $1.2 \mu\text{m}$ ).

The CSA and CSA-Shaper are continuously sensitive and require no external components or adjustment to be operated. In its self-biased

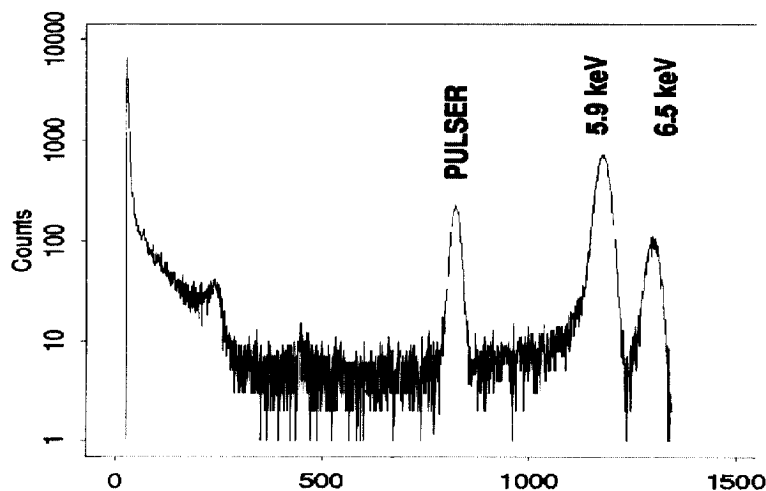


Fig. 8.  $^{55}\text{Fe}$  Spectrum of the AMS CSA at  $-70^\circ\text{C}$  (log vertical scale). FWHM = 111 eV (13e rms).

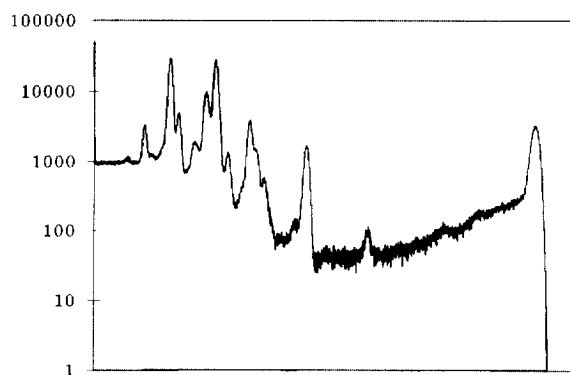


Fig. 9.  $^{241}\text{Am}$  Spectrum of the AMS CSA at  $-70^\circ\text{C}$  (log vertical scale).

condition, the AMS version has  $1.8\text{ G}\Omega$  feedback resistance. The linearity of the CSA voltage–charge gain has been improved by using a T-Network for the DC feedback (within 0.1% up to 1.8 fC input charge).

The CSA-Shaper uses an Adaptive Pole-Zero cancellation technique to accurately cancel the CSA Pole.

The best noise performance is obtained with MF pinched off, no detector connected, and with the circuit cooled to  $-70^\circ\text{C}$ . Under these conditions the ENC is 8e rms at a shaping time of  $8\text{ }\mu\text{s}$ , limited entirely by the flicker noise of the input transistor. This increases to 14e rms if we allow MF to be biased by the self-adaptive circuit. With a cooled SDD connected, a FWHM of 111 eV is obtained at  $2.4\text{ }\mu\text{s}$  corresponding to an ENC of 13e. The noise performance is the best ever achieved with a CMOS preamplifier.

Table 1

Noise performance of the CSA for various operating conditions

ENC (e rms)	$T_s$ ( $\mu\text{s}$ )	Detector	Cooled	Pinched
19	1.2	No	No	No
14	1.2	No	Yes	No
9	12	No	No	Yes
8	8	No	Yes	Yes
13	2.4	Yes	Yes	Yes

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